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Chemical Physics Letters 391 (2004) 334-337

CHEMICAL PHYSICS LETTERS

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# The Cl to NCl branching ratio in 248-nm photolysis of chlorine azide

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Received 2 April 2004 Available online 1 June 2004

# Abstract

The primary reaction products from 248-nm chlorine azide photolysis are identified in a collision-free experiment. In contrast to all previous reports, the radical channel producing  $Cl + N_3$  (95 ± 3%) is seen to dominate the photochemistry. The molecular channel producing  $NCl + N_2$  (5 ± 3%) was also observed.

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#### 1. Introduction

This Letter reports the first direct measurement of the product branching ratio from photolysis of ClN<sub>3</sub>. The photochemistry of this species [1] is of considerable interest because one of its photoproducts, NCl( $\tilde{a}$  <sup>1</sup> $\Delta$ ), the isovalent analog of singlet-oxygen, is an important energy carrier in an all gas-phase iodine laser [2,3] (AGIL). While the molecular mechanism of the AGIL is not completely understood, the near-resonant energy transfer reaction

$$\operatorname{NCl}(\tilde{a}^{1}\Delta) + \operatorname{I}(^{2}P_{3/2}) \to \operatorname{NCl}(\tilde{X}^{3}\Sigma) + \operatorname{I}(^{2}P_{1/2})$$
(1)

and slow relaxation of NCl( $\tilde{a}^{1}\Delta$ ) are important factors in its performance. The first reported observation of iodine lasing induced by NCl( $\tilde{a}^{1}\Delta$ ) employed 193-nm photolytic production of NCl( $\tilde{a}^{1}\Delta$ ) in mixtures of ClN<sub>3</sub> and CH<sub>2</sub>I<sub>2</sub> [4]. In 1998, gain in a dc-discharge-driven chemical laser was reported [5]. The most recent prototype (AGIL-2) can produce more than 15 W of continuous radiation at  $1.315 \mu m$  [6]. Such a light source has important potential applications in defense-related and medical research.

The importance of understanding the primary molecular processes relevant to the AGIL has led to a large number of laser-based kinetics studies, which utilize ClN<sub>3</sub> photolysis as a source of NCl( $\tilde{a}^1\Delta$ ) [7–11]. Despite intensive effort, the fundamental photochemistry of ClN<sub>3</sub> is not yet understood. For example, only limited data is available concerning the relative importance and thermochemistry of the dissociation pathways, as listed below:

 $ClN_{3}(\tilde{X}^{1}A') + h\nu \rightarrow Cl(^{2}P_{J}) + N_{3}(\tilde{X}^{2}\Pi, cyclic)$ (2a)

 $\rightarrow \operatorname{Cl}({}^{2}P_{J}) + \operatorname{N}_{2}(\tilde{X}^{1}\Sigma_{g}^{+}) + \operatorname{N}({}^{2}D_{J}, {}^{4}S_{3/2}) \quad (2b)$ 

$$\rightarrow \mathrm{NCl}(\mathrm{X}^{3}\Sigma, \tilde{\mathrm{a}}^{1}\Delta, \mathrm{b}^{1}\Sigma) + \mathrm{N}_{2}(\mathrm{X}^{1}\Sigma_{\mathrm{g}}^{+})$$
 (3a)

The above kinetics studies and earlier photodissociation studies [1,12] of ClN<sub>3</sub> point to NCl( $\tilde{a}^{1}\Delta$ ) (from reaction 3a) as the dominant photochemical product. However, none explicitly examine the Cl+N<sub>3</sub> channels, and all are vulnerable to problems related to the high reactivity of the primary products. For example, the photochemically-driven AGIL exhibited an induction time to reach the lasing threshold that could be modeled by a reaction mechanism where NCl( $\tilde{a}^{1}\Delta$ ) is

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<sup>0009-2614/\$ -</sup> see front matter  $\odot$  2004 Elsevier B.V. All rights reserved. doi:10.1016/j.cplett.2004.04.113

produced by chain reaction decomposition of ClN<sub>3</sub> [4]. Experiments employing 248-nm photolysis of ClN<sub>3</sub> also suggested that NCl( $\tilde{a}^{1}\Delta$ ) is produced by secondary reactions [8]. On the other hand, a recent study presented kinetic data that appeared to reflect the prompt production of NCl( $\tilde{a}^{1}\Delta$ ) as a primary product of ClN<sub>3</sub> photolysis [11]. Thus, there remains significant confusion regarding even the identity of the primary reaction products resulting from UV photolysis of ClN<sub>3</sub>.

Only recently have collision-free studies of ClN<sub>3</sub> photochemistry been reported [13–15]. In these studies, both primary Cl(<sup>2</sup>P<sub>J</sub>) [13] and N<sub>2</sub>( $\tilde{X}^{1}\Sigma_{g}^{+}$ ) [15] photoproducts were detected. The NCl( $\tilde{a}^{1}\Delta$ )/NCl ( $\tilde{X}^{3}\Sigma$ ) branching (reaction 3a) at 203-nm was also reported [15]. An unusual bimodal translational energy distribution of primary Cl-atom product [13] was consistent with photochemical production of a cyclic form of N<sub>3</sub> that had been predicted by theory [16]. Finally, a comprehensive and critical re-evaluation of the thermochemistry of all molecules containing a single Cl atom and three N atoms was put forth [13,14].

Here, we report results on the collision-free, 248-nm photolysis of ClN<sub>3</sub> using molecular beam photofragment translational spectroscopy [17,18]. This experiment, in addition to eliminating complications from secondary processes, enables one to observe all channels on an approximately equal footing. In contrast to all previous reports, we find clear evidence that Cl-atom production (2a) is by far the dominant (95  $\pm$  3%) collision-free dissociation channel. Channel (2b) occurs only through secondary collision-free decomposition of highly energetic N<sub>3</sub>. NCl production (3a) is an observed minor channel (~5  $\pm$  3%).

## 2. Experimental

The molecular beam photodissociation instrument on which these experiments were performed has been described in detail previously [19,20]. ClN<sub>3</sub> was formed by passing a mixture of 5% Cl<sub>2</sub> in He over the surface of moist sodium azide (NaN<sub>3</sub>). A standard drying agent (Drierite<sup>®</sup>) was used to remove water from the ClN<sub>3</sub>/ Cl<sub>2</sub>/He mixture at the exit of the generator. Under normal operating conditions, mass spectral analysis of the beam showed NCl<sup>+</sup>, an ionizer-induced fragment of ClN<sub>3</sub><sup>+</sup> [14], and little or no Cl<sub>2</sub><sup>+</sup>. Under conditions where Cl<sub>2</sub><sup>+</sup> signal could be detected, the NaN<sub>3</sub> sample was replaced. Using a backing pressure of approximately 0.5 bar, the mixture was formed into a pulsed molecular beam by expanding it through a piezoelectrically-actuated pulsed valve [21].

After passing through two collimating skimmers, the molecular beam intersected a 28-ns pulsed ultraviolet beam of unpolarized light produced by a LambdaPhysik<sup>®</sup> KrF laser (248 nm) operating at 100 Hz with 72 mJ pulse energy and focused to a  $\sim 10 \text{ mm}^2$  spot at the region of intersection. Photofragments scattered in the plane of the molecular and laser beams were detected as a function of laboratory scattering angle  $\Theta$  with a rotatable, triply differentially pumped mass spectrometer; after passing through three collimating apertures, the photofragments were ionized by electron impact  $(\sim 160 \text{ eV})$ , mass-selected with a quadrupole mass filter, and detected with a Daly-style detector [22]. At various scattering angles, photofragment time-of-flight (TOF) distributions were taken for ions at selected massto-charge ratios  $(m/e^{\circ}s)$  using a multi-channel scaler interfaced to a PC. The velocity distribution of the parent ClN<sub>3</sub> molecular beam was characterized by TOF at  $\Theta = 0^{\circ}$  using a retractable slotted chopper wheel to modulate the molecular beam.

## 3. Results and analysis

Fig. 1 shows TOF distributions for m/e = 35 (Cl<sup>+</sup>) and 49 (NCl<sup>+</sup>) at several scattering angles,  $\Theta$ . The data are represented by open circles. The solid lines are the results of forward-convolution simulations of the experiment using optimized center-of-mass translational energy  $(P(E_T))$  distributions shown in Fig. 2 for the  $Cl + N_3$  and  $NCl + N_2$  channels. Other input data required for analysis include the molecular beam velocity distribution, the finite size of the intersection and ionization regions, the finite angular acceptance angle of the detector and the laser polarization [23]. The  $P(E_T)$ distributions can also be used to generate TOF distributions of the momentum-matched counter-fragments (i.e. N<sub>3</sub> for Cl and N<sub>2</sub> for NCl), yielding distributions (not shown) that match components of the experimental TOF spectra at m/e = 42 (N<sub>3</sub><sup>+</sup>) and 28 (N<sub>2</sub><sup>+</sup>), although the latter is complicated by secondary dissociation of the N<sub>3</sub> primary product.

Successfully fitting the data allowed us to derive the center-of-mass frame flux of Cl and NCl reaction products after proper correction for the ionization efficiency of each species [24]. This procedure was simplified by the fact that electron impact induced fragmentation of NCl and Cl was insignificant. The relative total ionization efficiency was estimated by using a known empirical correlation between molecular polarizability and total electron ionization cross-section, which were calculated using the additivity rule proposed by Fitch and Sauter [25] as explained by Schmoltner [26]. Briefly, the electron impact cross section (in units of  $10^{-16}$  cm<sup>2</sup>) is estimated from the following additivity formula:

$$Q = 0.082 + \sum_{i} a_{i} n_{i}.$$
 (4)



Fig. 1. TOF spectra for m/e = 35 (Cl<sup>+</sup>) and 49 (NCl<sup>+</sup>) at detection angles ( $\Theta$ ) of 10°, 30° and 50°. For all TOF spectra the open circles represent the data and the solid line represents the forward convolution fit to the data using the P( $E_T$ ) distributions in Fig. 2. Note the appearance of a 'fast' (dashed) and a 'slow' (dotted) channel in the m/e = 35 TOF spectra.

This analysis resulted in a branching fraction:

$$\phi_{\rm Cl} \equiv \frac{[\rm Cl]}{[\rm NCl] + [\rm Cl]} = 0.95 \pm 0.03.$$
 (5)

This result stands in marked contrast to the 'conventional wisdom' that the Cl + N<sub>3</sub> channel is relatively minor. This assumption has been inferred from observations of the NCl+N2 channel alone, not by any systematic comparison of the two channels. For example, measurements by Henshaw et al. [9] of the NCl( $a^{1}\Delta$ ) concentration at early times via titration with Br2 gave a quantum yield of 60% for NCl( $a^{1}\Delta$ ) + N<sub>2</sub> at 193 nm. Recent work by Komissarov et al. [11] which reported a rapid appearance of NCl( $\tilde{a}^1 \Delta$ ), also appeared to provide strong evidence for the dominance of the  $NCl + Cl_2$ channel. However, analyses of these experiments are based on the assumption that the NCl is a primary product, not the result of rapid secondary reactions. Heaven and co-workers [27] have since obtained much improved kinetic results which strongly indicate that NCl( $\tilde{a}^1\Delta$ ) is indeed formed by secondary reactions and the mechanism for its production under their conditions has been elucidated. Our result is quite unambiguous, since our experiment is conducted under collisionless conditions and the two channels were investigated using the same type of measurement.

It is not possible to identify the electronic states of NCl due to the fact that a large amount of the available energy appears as product rotation and/or vibration in the molecular channel. We note that the bimodal  $P(E_T)$  distribution for the Cl + N<sub>3</sub> channel needed to fit the two peaks in the TOF spectra is quite similar to that obtained by Hansen et al. [13] in a recent photofragment velocity-map imaging experiment, in which the slower feature was interpreted as possible evidence for the photolytic production of cyclic-N<sub>3</sub>. Thus, the experiments reported in this work confirm the existence of a bimodal translational energy distribution in the radical channel. This means that N<sub>3</sub> is formed in a low energy (linear azide radical) and a high energy (possibly cyclic-



Fig. 2. Center-of-mass translational energy distributions,  $P(E_T)$ , for the radical (a) and molecular (b) channels in ClN<sub>3</sub> photodissociation. The arrows indicate the maximum allowed translation energy release when products have no rotational or vibrational excitation. The distributions shown are able to produce the fit to all of the TOF data shown in Fig. 1.

 $N_3$ ) form. The branching fraction for the radical channel found here suggests that the high energy form of  $N_3$ formed in ClN<sub>3</sub> photolysis is produced with large quantum efficiency.

#### Acknowledgements

The authors thank Prof. Michael Heaven for reviewing an early version of this manuscript. AMW acknowledges the support of an Air Force Office of Scientific Research contract under grant number F49620-95-1-0234. NH acknowledges the support of the Alexander von Humboldt Foundation under a Fyodor-Lynen Stipend. DMN acknowledges support from the Air Force Office of Scientific Research under Grant No. F49620-03-1-0085.

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